trans-bent configuration, we point out that the former has a less repulsive alignment of the  $F_2C=CF$  group dipole moments than does the latter. However, the final answer to this question is bound to be much more complicated.

Acknowledgments. It is a pleasure to thank Dr. S. L. Manatt for a sample of 1,1,4,4-tetrafluorobutadiene, and Mr. N. A. Kuebler for running the optical spectra for us. Additionally, discussions with Professor V. W. Laurie and Dr. A. Andreassen are acknowledged with thanks.

## Ground States of Conjugated Molecules. XXII. Polarographic Reduction Potentials of Hydrocarbons<sup>1</sup>

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Abstract: The polarographic reduction potentials of 76 aromatic hydrocarbons correlate well with electron affinities calculated by a recently described semiempirical SCF MO procedure. Points of attachment of the first two hydrogen atoms are predicted; in the few cases where data are available they confirm the predictions. The available values for polarographic oxidation potentials correlate well with calculated ionization potentials.

Correlations between the polarographic reduction half-wave potentials  $(-E_{1/2})$  of conjugated hydrocarbons and the lowest unoccupied HMO orbital energies  $(-m_{n+1})$  have been described by many authors, <sup>3-7</sup> giving in all cases reasonably good correlations with experiment, and some analogous correlations have also been reported using the Pariser-Parr method.<sup>8</sup> Here we present a correlation between the polarographic reduction half-wave potentials and electron affinities calculated by the variable- $\beta$  SCF MO method recently developed in this laboratory.<sup>9-12</sup> We have also used this procedure to calculate the points of attachment of the first two hydrogen atoms during reduction by electron transfer, and finally, the available data concerning polarographic oxidation potentials have been correlated with calculated ionization potentials.

## **Theoretical Method**

Reduction of an aromatic hydrocarbon (R) by a reversible one-electron transfer leads to the radical anion  $(R^{-})$ 

$$\mathbf{R} + \mathbf{e} \rightleftharpoons \mathbf{R}^{-} \tag{1}$$

- (3) A. Maccoll, Nature (London), 163, 178 (1948).
- (4) L. E. Lyons, ibid., 166, 193 (1950).
- (5) A. Pullman, B. Pullman, and G. Berthier, Bull. Soc. Chim. Fr., (1950).
  (6) G. J. Hoijtink and J. van Schooten, Recl. Trav. Chim. Pays-Bas, 71,
- 1089 (1952).
  - (7) G. J. Hoijtink, ibid., 75, 1525 (1955).
  - (8) J. Koutecký, Z. Phys. Chem. (Frankfurt am Main), 52, 8 (1967).
- (9) A. L. H. Chung and M. J. S. Dewar, J. Chem. Phys., 42, 756 (1965).
- (10) M. J. S. Dewar and G. J. Gleicher, J. Amer. Chem. Soc., 87, 685 (1965).
  - (11) M. J. S. Dewar and C. de Llano, ibid., 91, 789 (1969).

The free energy change  $(\Delta F^{\circ})$  for the above process in solution is as follows

$$\Delta F^{\circ} = (F^{\circ}_{R})_{gas} - (F^{\circ}_{R})_{gas} + (F^{\circ}_{electron})_{Hg} + (\Delta F^{\circ})_{solv} \quad (2)$$

The difference  $(F^{\circ}_{R})_{gas} - (F^{\circ}_{R})_{gas}$  is equal to the electron affinity (A) of the molecule R in the gas phase, if we neglect entropy effects.

On the other hand, one can write down the halfwave potential  $(E_{1/2}^{red})$  for a reversible one-electron electrode process as follows

$$-E_{1/2}^{\text{red}} = \frac{\Delta F^{\circ}}{\mathfrak{F}} - \frac{RT}{\mathfrak{F}} \ln \left(\frac{D_{\text{R}}}{D_{\text{R}}}\right)^{1/2} + \frac{RT}{\mathfrak{F}} \ln \frac{F_{\text{R}}}{F_{\text{R}}} \quad (3)$$

where  $\mathcal{F}$  is the Faraday constant, R is the gas constant, T is temperature in degrees Kelvin, and  $D_{\rm R}$ and  $D_{R^-}$  and  $F_R$  and  $F_{R^-}$ , are, respectively, the diffusion coefficients and activity coefficients of R and R-. Since these should be very similar, terms involving them may reasonably be neglected; e.g., (3) then becomes

$$-E_{1/2}^{\text{red}} = \frac{\Delta F^{\circ}}{\mathfrak{F}}$$
(4)

Substituting A into eq 2 using eq 4

$$-E_{1/2}^{\text{red}} \frac{1}{\Im} [A + (F^{\circ}_{\text{electron}})_{\text{Hg}} + (\Delta F^{\circ})_{\text{solv}}] \quad (5)$$

The last two terms in eq 5, for polarographic reduction of a series of similar compounds under similar experimental conditions, should be constant. In this case

$$-E_{1/2}^{\text{red}} = \frac{A}{\mathfrak{F}} + \text{ constant}$$
 (6)

A plot of  $-E^{1/2}$ <sup>red</sup> vs. A should then be linear.

A similar relationship should hold between the polarographic oxidation potentials of hydrocarbons  $(E_{1/2}^{\text{ox}})$  and calculated ionization potentials (I); *i.e.* 

$$E_{1/2}^{\text{ox}} = I + \text{constant}$$
(7)

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<sup>(1)</sup> Part XXI: M. J. S. Dewar, A. J. Harget, N. Trinajstić, and S. D. Worley, *Tetrahedron*, in press. This work was supported by the Air Force Office of Scientific Research through Grant No. AF-AFOSR-70-1881.

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<sup>(12)</sup> M. J. S. Dewar and T. Morita, ibid., 91, 796 (1969).

Table I.	Calculated Heats of Atomization of Hydrocarbons and Hydrocarbon Radical Anions, Electron Affinities, Ob-	served
Polarogra	uphic Half-Wave Reduction Potentials, and the Lowest Unoccupied HMO Orbital Energies	

	Heat of atomization, eV					Heat of atomization, eV					
Molecule	Hydro- carbon	Radical anion	A, eV	$-E_{1/2}^{\mathrm{red}},$ $\mathrm{V}^{a}$	$-m_{n+1}{}^{b}$	Molecule	Hydro- carbon	Radical anion	A, eV	$-E_{1/2}^{\mathrm{red}},$ V <sup>a</sup>	$-m_{n+1}^{b}$
I	90.612	81.025	0.074	1.98	0.6180°	XXXIX	201.527	192, 506	0.640	1.64	0.5392°
II	123.300	114.893	0.653	1.46	0.4142°	XL	239.390	230.776	1.046	1.28	0.3557i
III	157.112	148.512	1.060	1.14	0.2950°	XLI	249.232	240.960	1.388	0.92	0.2539/
IV	190. <b>29</b> 4	181.975	1.341	0.86	0.2197°	XLII	287.091	279.095	1.665	0.68	0.1826 <sup>i</sup>
v	124.222	114.835	0.273	1.94	0.6052°	XLIII	287.099	279.042	1.603	0.70	0.1916
VI	157.580	148.560	0.640	1.53	0.4523°	XLIV	138.627	129.631	0.664	1.61	0.4450°
VII	190.853	182.196	1.003	1.19	0.3271ª	XLV	172.380	164.330	0.610	1.67	0,4 <b>970</b> ⁰
VIII	224.031	215.663	1.292	0.95	0.2436°	XLVI	172.111	163.381	0.930	1.36	0.3711
IX	157.943	148.535	0.251	1.97	0.6840°	XLVII	205,810	197.077	0.926	1.08	0.3983/
X	191.348	182.257	0,568	1.54	0.4991	XLVIII	205.902	197.093	0.851	1.45	0.42161
XI	224,636	215.923	0.948	1.21	0.35571	XLIX	206.114	197.015	0.562	1.69	0.5550 <sup>d</sup>
XII	257.825	249.420	1.254	0.93	0.2621	L	205.533	197.057	1.185	1.16	0.30274
XIII	191.238	182,232	0.654	1.55	0.4735°	LI	205,646	197.043	1.057	1.32	0.3420 <sup>d</sup>
XIV	224.555	215.851	0.956	1.25	0.3584/	LII	205.805	196.664	0.519	1.51	0.5053*
XV	257.760	249.341	1.242	0.85	0.2691	LIII	205.499	196.987	1.148	1.15	0.3202*
XVI	157.772	148.628	0.516	1.81	0.5201°	LIV	272.233	264.088	1.515	0.78	0.2130
XVII	191, 113	182, 253	0.800	1.44	0.4048/	LV	272,483	263.978	1.155	1.15	0.2734
XVIII	224,431	215.735	0.964	1.33	0.3482/	LVI	239,461	230.762	0.961	1.35	0.3957
XIX	157.781	148.505	0.384	1.75	0.5676	LVII	220,270	211.856	1.247	1.18	$0.2846^{d}$
XX	191,238	182.161	0.583	1.57	0.4917	LVIII	287.749	279.279	1.190	1.15	0.3445i
XXI	191 458	182 398	0.600	1 65	0.5319/	LIX	186.431	177.951	1.180	1.19	0.2910 <sup>d</sup>
XXII	224,982	215,927	0.605	1.57	0.5224/	LX	253.537	245, 261	1.384	1.06	0.2580*
XXIII	225,108	216.229	0.781	1.59	0.5115/	LXI	253.177	245,153	1.636	0.80	0.1891
XXIV	191 361	182 201	0.500	1 73	0.5498/	LXII	138.668	129.848	0.840	1.35	0.3709
XXV	191, 113	182,210	0.757	1.40	0.4186/	LXIII	171.987	163.558	1.231	0.98	0.25214
XXVI	191 361	182, 292	0.591	1 79	0.5019	LXIV	205.236	197.178	1.603	0.65	$0.1604^{d}$
XXVII	190.980	182.027	0.707	1.53	0.4372°	LXV	172.279	163.555	0.935	1.38	0.3770 <sup>d</sup>
XXVIII	224 263	215 577	0.974	1 22	0.3357*	LXVI	172,148	163.512	1.024	1.17	0.3116 <sup>d</sup>
XXIX	224 618	215 709	0.751	1 50	0.4287	LXVII	172,118	163.291	0.834	1.39	0.40074
XXX	258 263	249 382	0 779	1 52	0.4321	LXVIII	205.445	196.930	1.146	1.06	0.2666
XXXI	257 991	249 129	0 798	1.47	0.3940°	LXIX	205,405	196.597	0.852	1.38	0.3966 <sup>h</sup>
XXXII	172 146	163 442	0.956	1 25	0.3473	LXX	186.710	178,420	1.371	0.90	0.18594
XXXIII	205 488	197 076	1 248	0.97	0.2648/	LXXI	220.050	212 092	1.702	0.60	0.08804
XXXIV	238 805	230 607	1 463	0.88	0.2135/	LXXII	253.322	245.702	2.040	0.34	0.00004
XXXV	239 095	230 731	1 296	1 00	0.2673	LXXIII	109.753	100 156	0.063	2.08	0.7046
XXXVI	186 893	178 011	0 779	1 49	0.4392°	LXXIV	162 352	153 106	0.414	1.91	0.5926
XXXVII	220 594	211 749	0.815	1 52	0.4631	LXXV	214.951	205 859	0.568	1.77	0.5361
XXXVIII	254.131	245.432	0.961	1.35	0.4114/	LXXVI	205.221	196.967	1.406	0.46	0.19874

<sup>a</sup> The experimental half-wave potentials are taken from I. Bergman, *Trans. Faraday Soc.*, **50**, 829 (1954). <sup>b</sup>  $\alpha + m_{n+1}\beta$  is the energy of the lowest unoccupied HMO orbital. <sup>c</sup>C. A. Coulson and A. Streitwieser, Jr., "Dictionary of  $\pi$ -Electron Calculations," Pergamon Press, Oxford, 1965. <sup>d</sup> A. Streitwieser, Jr., and J. I. Brauman, "Supplemental Tables of Molecular Orbital Calculations," Vol. I and II, Pergamon Press, Oxford, 1965. <sup>e</sup> P. Hochmann, J. Dubský, J. Koutecký, and C. Párkányi, *Coll. Czech. Chem. Comm.*, **30**, 3560 (1965). <sup>f</sup> R. Zahradník and C. Párkányi, *ibid.*, **30**, 3536 (1965). <sup>e</sup> M. Titz and P. Hochmann, *ibid.*, **31**, 4168 (1966). <sup>h</sup> P. Hochmann, J. Koutecký, and R. Zahradník, *ibid.*, **27**, 3053 (1962). <sup>i</sup> R. Zahradník, J. Michl, and J. Koutecký, *ibid.*, **29**, 3184 (1964). <sup>i</sup> Calculated by us.

The reduction of a hydrocarbon R to its dihydro derivative  $RH_2$ , by electron transfer in the presence of a source of protons, normally takes place in two steps *via* an intermediate radical  $RH_{\cdot}$ ; *viz*.

$$R + e \longrightarrow R^{-}$$

$$R^{-} + H^{+} \longrightarrow RH^{-}$$

$$RH^{-} + e \longrightarrow RH^{-}$$

$$RH^{-} + H^{+} \longrightarrow RH_{2}$$
(8)

Since protonation of a mesomeric carbanion is an extremely facile process, it should take place preferentially at the position with maximum formal negative charge.<sup>13</sup> The point of attachment of the first proton should then be the position with maximum  $\pi$ -electron density in the ion radical R<sup>-</sup>, and the point of attachment of the second proton the corresponding position in the ion RH<sup>-</sup>.

(13) (a) M. J. S. Dewar, *Discuss. Faraday Soc.*, 2, 261 (1947); (b) "The Molecular Orbital Theory of Organic Chemistry," McGraw-Hill, New York, N. Y., 1969.

The electron affinities (A) and ionization potentials (I) of hydrocarbons (R) were calculated from the following expressions

$$A = \Delta H_{a}(\mathbf{R}^{-}) - \Delta H_{a}(\mathbf{R})$$
(9)

$$I = \Delta H_{a}(\mathbf{R}) - \Delta H_{a}(\mathbf{R}^{+})$$
(10)

where the  $\Delta H_a$ 's are heats of atomization calculated by our SCF MO  $\pi$  approximation,<sup>11</sup> those for the ion radicals R<sup>+</sup> and R<sup>-</sup> being found by the "half-electron" variant.<sup>14</sup>

## Results

In Table I are shown heats of atomization calculated for the 76 aromatic hydrocarbons listed in Figure 1, and for the radical anions formed by their reduction; some of these values had already been reported in part IX of this series.<sup>14</sup> In comparing the corresponding electron affinities (third column of Ta-

(14) M. J. S. Dewar, J. A. Hashmall, and C. G. Venier, J. Amer. Chem. Soc., 90, 1953 (1968).



Figure 1. Geometries of molecules studied. Circles denote the predicted point of addition of the first proton and squares that of the second proton.

ble I) with experiment, it is necessary to use reduction potentials measured under identical experimental conditions, for although the relative values for different hydrocarbons seem to be the same in different solvents, etc.,<sup>15</sup> their absolute values vary. We have therefore confined ourselves to the values reported by Bergman,<sup>16</sup> since no other measurements refer to so large a number of hydrocarbons. His values for  $E_{1/2}^{\text{red}}$  are shown in the fourth column of Table I, while Figure 2 shows a plot of them *vs.* our calculated electron affinities. It will be seen that there is an excellent linear relation between the electron affinities and reduction potentials, the correlation coefficient being 0.97 and the standard deviation only 0.028 V. The slope of the least-squares line (0.92) is slightly less than the theoretical value of unity (eq 6), but the difference is within the limits of experimental error<sup>16</sup> in the measured values of  $E^{1/2}$ <sup>red</sup>. The calculated electron affinities also agree closely with experiment.<sup>14,17</sup>

The fact that the slope of the line in Figure 2 is so close to unity must mean that the differences in solva-(17) R. S. Becker and E. Chen, J. Chem. Phys., 45, 2409 (1966).

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<sup>(15)</sup> P. A. Given, J. Chem. Soc., 2684 (1958).

<sup>(16)</sup> I. Bergman, Trans. Faraday Soc., 50, 829 (1954).

tion energy between the hydrocarbons and the corresponding radical anions either vary little or are small; the latter seems to be more likely in view of the insensitivity of the relative values for different hydrocarbons to changes in the solvent,<sup>15</sup> and because the aprotic solvents used in these studies should not solvate anions at all efficiently. However, the small deviation in the slope of the line in Figure 2 from the theoretical value



Figure 2. Plot of electron affinities (A) vs. polarographic half-wave reduction potentials  $(E_{1/2}^{\text{red}})$  for the compounds listed in Figure 1.

of unity may reflect small changes in solvation effects along the series.

As we remarked earlier, there is also a good linear relation between  $E^{1/2^{red}}$  and the energies of the lowest unoccupied Hückel MO's;<sup>3-7,18</sup> in order to compare this with the one indicated in Figure 2, we had to carry out HMO calculations for many of the hydrocarbons since these were not listed in the available compilations.<sup>19</sup> Energies of the corresponding lowest unoccupied MOs are shown in the fifth column of Table I. The correlation between them and the observed reduction potentials is quite good (correlation coefficient, 0.97) but the standard deviation (0.077 V) is almost three times ours; moreover, the HMO results are expressed in terms of a parameter  $(\beta)$ , whereas ours are absolute values, in electronvolts. This difference is quite an important one, for deviations of the slopes of plots such as that of Figure 2 from the theoretical value provide a good indication of any interference due to solvent effects. And finally, the success of HMO theory in this connection is almost certainly due to the fact that nearly all the compounds studied here are alternant hydrocarbons; it is well known<sup>13b</sup> that the HMO method is reasonably reliable only for this particular class of compound, whereas the

(18) A. Streitwieser, "Molecular Orbital Theory for Organic Chemists," Wiley, New York, N. Y., 1961, pp 173-185.

(19) (a) C. A. Coulson and A. Streitwieser, Jr., "Dictionary of  $\pi$ -Electron Calculations," Pergamon Press, Oxford, 1965; (b) A. Streitwieser, Jr., and J. I. Brauman, "Supplemental Tables of Molecular Orbital Calculations," Vol. I and II, Pergamon Press, Oxford, 1965; (c) P. Hochmann, J. Dubský, J. Koutecký, and C. Párkányi, *Coll. Czech. Chem. Commun.*, **30**, 3560 (1965); (d) R. Zahradník and C. Párkányi, *ibid.*, 30, 3536 (1965); (e) M. Titz and P. Hochmann, *ibid.*, 31, 4168 (1966); (f) P. Hochmann, J. Koutecký, and R. Zahradník, *ibid.*, 27, 3053 (1962); (g) R. Zahradník, J. Michl, and J. Koutecký, *ibid.*, 29, 3184 (1964).

procedure used here should be equally applicable to compounds containing heteroatoms. Indeed, a good example of both these points is provided by a recent study of reduction potentials of quinones.<sup>20</sup> Here again there is a good linear relation between the observed reduction potentials and electron affinities calculated by our SCF MO method, but the slope of the line (0.39) was much less than the theoretical value (0.5); here the reductions are carried out in protic solvents where the anions are very strongly solvated.

As indicated earlier, we also studied the reduction of the hydrocarbons to dihydro derivatives by the process indicated in eq 8. The first hydrogen should attach itself to the atom with the highest  $\pi$ -electron density in the radical anion  $R^-$ ; these positions are indicated by circles in Figure 1. The second hydrogen should likewise attach itself to the position with maximum  $\pi$ -electron density in the intermediate anion RH<sup>-</sup>; this is indicated in Figure 1 by a square. In some cases two or more positions in RH- have essentially identical  $\pi$ -electron densities; all of them are then marked in Figure 1. Unfortunately, few of the reduction products of these hydrocarbons have been studied; indeed, data seem to be available only for benzene and naphthalene. The former reduces 1,4 exclusively, as we predict (the well-known Birch reduction<sup>21</sup>), and the same is also true of naphthalene.<sup>22</sup> In both cases the HMO method predicts equal ease of 1,2 and 1,4 addition. There are a number of other similar discrepancies between our predictions and those<sup>6</sup> of HMO theory; we hope that our calculations may stimulate further experimental work in this area.

The polarographic oxidation potentials of hydrocarbons should be related to their ionization potentials in the same way that the reduction potentials are to electron affinities (eq 6). Polarographic oxidation potentials have been measured by several groups of workers,<sup>23-26</sup> and the results correlate quite well with the energies of the highest occupied Hückel MO's for the hydrocarbons.<sup>27</sup> Gleicher and Gleicher<sup>28</sup> have also applied an earlier version<sup>10</sup> of the present SCF MO treatment to this problem; they found a good correlation between their calculated orbital energies and the oxidation potentials of Pysh and Yang.<sup>25</sup>

Table II shows ionization potentials calculated by Koopmans' theorem, which has been shown<sup>14</sup> to give results very similar to those from direct calculations using the "half-electron" method, <sup>14</sup> together with the polarographic oxidation potentials ( $E_{1/4}^{\circ x}$ ) of Pysh and Yang;<sup>25</sup> these were chosen since they represent the most complete series of measurements made under identical experimental conditions. As will be seen from Figure 3, there is an excellent linear relation between the calculated ionization potentials and the measured

(20) M. J. S. Dewar and N. Trinajstić, *Tetrahedron*, 25, 4529 (1969).
(21) (a) A. J. Birch, *Quart. Rev., Chem. Soc.*, 4, 69 (1950); (b) ref 13b, p 292.

(22) S. Wawzonek and H. A. Laitinen, J. Amer. Chem. Soc., 64, 2365 (1942).

(23) H. Lund, Acta Chem. Scand., 11, 1323 (1957).

- (24) J. W. Loveland and G. R. Dimeler, Anal. Chem., 33, 1196
  (1961).
  (25) E. S. Pysh and N. C. Yang, J. Amer. Chem. Soc., 85, 2124
- (1963). (26) W. C. Neikam, G. R. Dimeler, and M. M. Desmond, J. Electro-
- chem. Soc., 111, 1190 (1964).
  (27) G. J. Hoijtink, Recl. Trav. Chim. Pays-Bas, 77, 555 (1955).

(28) G. J. Gleicher and M. K. Gleicher, J. Phys. Chem., 71, 3693 (1967).

 Table II.
 Calculated Vertical Ionization Potentials,

 Observed Polarographic Half-Wave Oxidation Potentials,
 and the Highest Occupied HMO Orbital Energies

Molecule	-(Vertical ionization potential), eV	$-E_{1/2^{\mathrm{ox}}},$ $\mathrm{V}^a$	$+m_n^b$
Bangana			
(D)	0 369	2 30	1 0000
(U) <sup>2</sup>	9.500	2.50	0.6180
I TT	0.341	1.04	0.0100
11	7.512	0.77	0.4142
	7.597	1.50	0.2930
V VT	0.434	1.50	0.0032
	0.037	1,10	0.4323
	0,434	1.55	0.0640
	0.090	1.23	0.4991
	8.070	1.19	0.4733
AVI VVII	0.193	1.35	0.5201
	1.8/5	1.01	0.4048
	8.095	1.20	0.4917
XXVI	8.157	1.33	0.5019
	7.724	0.85	0.3473
XXXVI	7.846	1.01	0.4392
XXXIX	7.976	1.23	0.5392
XLIV	7.951	1.16	0.4450
XLV	8.028	1.27	0.4970
XLVI	7.747	0.94	0.3711
XLVII	7.777	1.15	0.3983
XLVIII	7.841	1.01	0.4216
LXII	8.337	1.45	0.6180

<sup>a</sup> E. S. Pysh and N. C. Yang, J. Amer. Chem. Soc., 85, 2124 (1963). <sup>b</sup> HMO values from tables cited in Table I. <sup>c</sup> Polarographic half-wave reduction potentials were not measured owing to experimental difficulties; its reduction potential is too high.

 $E_{1/2}^{\text{ox}}$ , the correlation coefficient being 0.98 and the standard deviation only 0.035 V; indeed, the latter is comparable with the claimed experimental error.<sup>24</sup>

The slope of the line (0.87) in Figure 3 is, however, much less than unity; this presumably reflects the fact that carbonium ions are strongly solvated even



Figure 3. Plot of vertical ionization potentials (IP<sub>v</sub>) vs. polarographic half-wave oxidation potentials  $(E_{1/2}^{\text{ox}})$ .

in an aprotic solvent such as that (acetonitrile) used here. The more delocalized the charge in the radical cation formed by oxidation of a given hydrocarbon, the lower should be its energy; one would therefore expect a correlation between oxidation potential and energy of solvation of the radical cation formed by oxidation. Such a correlation would leave the relation between ionization potential and  $E_{1/2}^{\text{ox}}$  linear but reduce its slope; a similar reduced slope was observed<sup>20</sup> in a correlation of reduction potentials of quinones with calculated electron affinities and explained in the same way.

Here again the HMO method is clearly inferior. The best linear relation between  $E_{1/2}^{ox}$  and the energy of the highest occupied Hückel MO shows a standard deviation of 0.092 V, nearly three times ours, and the HMO method also suffers from the deficiencies pointed out above in the discussion of reduction potentials. The present procedure also gives a slightly better correlation than that of Gleicher and Gleicher,<sup>28</sup> as might be expected, since it is known to give even better estimates of heats of atomization of molecules.